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## Optimization of the base thickness of a silicon solar cell with multi-vertical junctions connected in series, through dynamic expressions of the recombination velocity on the back side

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**Abstract** The minority carrier's diffusion equation, is solved in the base of a vertical Multi- junction's silicon solar cell connected in series (VMJS), under illumination in frequency modulation. The AC photocurrent is studied as a function of the recombination velocity at the junction, and allow us to establish the expressions of the AC recombination velocity of minority carriers on the back side. The profile of AC recombination velocity of minority carriers leads to the determination of the optimum thickness of the base, allowing the production of a maximum of electric current. The optimum thickness is thus expressed by mathematical relations as a function of both, the dynamic diffusion coefficient and the frequency of modulation of the incident light.

**Keywords** VMJ Silicon Solar Cell – AC Recombination Velocity - Base Thickness – modulation frequency

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### 1. Introduction

The silicon solar cell (n+/p/p+) is produced under several architectures [1, 2, 3], according to both the intensity and the orientation of the incident illumination relative to the plane of the surface of the space charge region (n+/p) [4]. For incident illumination perpendicular to the plane of the junction, the solar cell may be monofacial [5] or bifacial for illumination carried out by the back side (p+) [4, 6, 7].

When the incident light arrives parallel to the plane of the junctions (n+/p) and (p/p+), the structure can be vertical multi-junctions connected in series (n+/p/p+) [8] or (n/p) connected in parallel [9]. It is therefore designed using electronically poor (Si) material under concentrated illumination intensity [10, 11], so that photogenerated charge carriers everywhere in the different regions can be easily collected, because of the proximity of the junctions (n/p) [12]. This constitutes an economic advantage.

This raises the problem of the distance to be covered by the minority carriers, in other words, it is therefore a second advantage to optimize, the thickness of the base that would produce a maximum photocurrent, by collecting more carriers.

The proposed work aims to determine the optimum thickness [13, 14, 15] of the base of the vertical multi-junctions silicon solar cell (n+/p/p+), connected in series, when it is under monochromatic illumination in frequency modulation.



For this, the diffusion equation in dynamic regime relative to the minority carriers in the base [16, 17] is solved, taking into account the recombination velocity ( $S_f$ ) at the junction and ( $S_b$ ) on the rear side (p/p+) [18, 19, 20, 21, 22, 23].

The expression of the dynamic photocurrent  $J_{ph}$  is established. It is represented graphically as a function of ( $S_f$ ). From this representation, the expressions of the dynamic recombination velocity ( $S_b$ ) on the back side (p/p+) are deduced [24] dependent on base depth

The analysis of these expressions in graphic mode, makes it possible to determine the optimum thickness ( $H_{opt}$ ) of the base, necessary for the production of an optimal photocurrent for given frequency. The mathematical modeling of this thickness as a function of both, the diffusion coefficient and the frequency of modulation of the illumination is established.

**2. Theory**

The structure of the  $n^+ - p - p^+$  vertical junction silicon solar cell [1, 7, 8] under monochromatic illumination, in frequency modulation is given by figure 1.

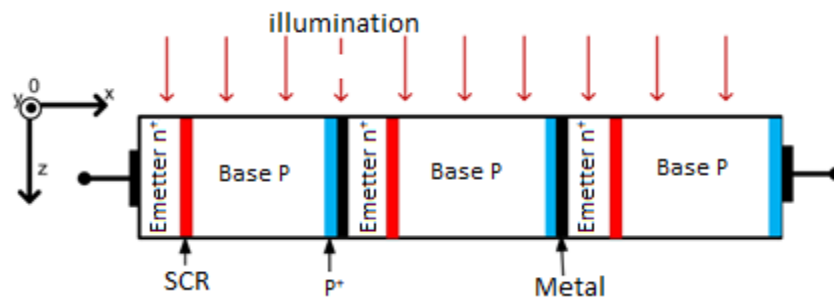


Figure 1: Structure of vertical junction solar cell under modulated monochromatic light

The series vertical junction solar cell, being a series combination of photovoltaic solar cells connected to each other by a metal, is designed in such a way that the incident illumination is parallel to the plane of the space charge region (SCR). The structure of this solar cell is shown in Figure 2:

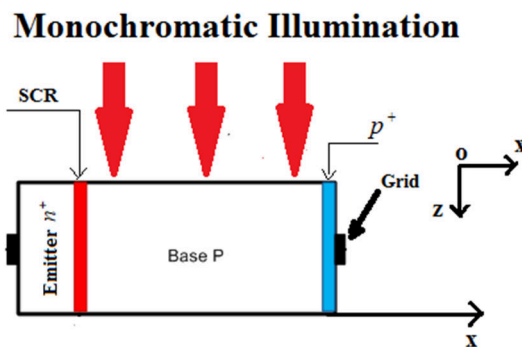


Figure 2: One unit of the series vertical junction solar cell

The excess minority carriers' density  $\delta(x, t)$  generated in the base of the solar cell obeying to the (2D) continuity equation [8, 25, 26, 27] is given as:

$$D(\omega) \times \frac{\partial^2 \delta(x, z, t)}{\partial x^2} - \frac{\delta(x, z, t)}{\tau} = -G(z, \omega, t, \alpha) + \frac{\partial \delta(x, z, t)}{\partial t} \tag{1}$$

The expression of the excess minority carriers' density is written, according to the space coordinates (x,z) and the time t, as:

$$\delta(x, z, t) = \delta(x, z) \cdot e^{-j\omega t} \quad (2)$$

Carrier generation rate  $G(z, \omega, t, \alpha)$  is given by the following relationship:

$$G(z, \omega, t, \alpha) = g(z, \alpha) \cdot e^{-j\omega t} \quad (3)$$

With:

$$g(z, \alpha) = \alpha(\lambda) \cdot I_0(\lambda) \cdot (1 - R(\lambda)) \cdot e^{-\alpha(\lambda)z} \quad (4)$$

- $z$  is the depth in the base.  $I_0$  is the intensity of monochromatic illumination. The quantities  $\alpha(\lambda)$  and  $R(\lambda)$  are the monochromatic absorption and reflection coefficients of the silicon material with wavelength ( $\lambda$ ) [28].
- $D(\omega)$  is the complex diffusion coefficient of excess minority carrier in the base. Its expression is given by the relationship [25, 27]:

$$D(\omega) = D_0 \times \left( \frac{1 - j \cdot \omega^2 \cdot \tau^2}{1 + (\omega\tau)^2} \right) \quad (5)$$

By replacing equations (2) and (3) in equation (1), the continuity equation for the excess minority carriers' density in the base is reduced to the following relationship:

$$\frac{\partial^2 \delta(x, z)}{\partial x^2} - \frac{\delta(x, z)}{L^2(\omega)} = -\frac{g(z)}{D(\omega)} \quad (6)$$

$L(\omega)$  is the complex diffusion length of excess minority carriers in the base given by:

$$L(\omega) = \sqrt{\frac{D(\omega)\tau}{1 + j\omega\tau}} \quad (7)$$

$\tau$  is the excess minority carriers lifetime in the base.

**The solution of equation (6) is:**

$$\delta(x, z, \omega, \alpha(\lambda)) = A \cdot \cosh\left[\frac{x}{L(\omega)}\right] + B \cdot \sinh\left[\frac{x}{L(\omega)}\right] + K \cdot e^{-\alpha \cdot z} \quad (8)$$

$$\text{With } K = \frac{\alpha \cdot I_0 \cdot (1 - R) \cdot [L(\omega)]^2}{D(\omega)[L(\omega)^2 \cdot \alpha^2 - 1]} \quad (9)$$

$$\text{and } (L(\omega)^2 \cdot \alpha^2 \neq 1) \quad (10)$$

Coefficients A and B are determined through the boundary conditions expressed as:

- At the junction ( $x = 0$ )

$$\left. \frac{\partial \delta(x, z, \omega, \alpha(\lambda))}{\partial x} \right|_{x=0} = S_f \cdot \left. \frac{\delta(x, z, \omega, \alpha(\lambda))}{D(\omega)} \right|_{x=0} \quad (11)$$

- On the back side in the base ( $x = H$ )

$$\left. \frac{\partial \delta(x, z, \omega, \alpha(\lambda))}{\partial x} \right|_{x=H} = -S_b \cdot \left. \frac{\delta(x, z, \omega, \alpha(\lambda))}{D(\omega)} \right|_{x=H} \quad (12)$$



Sf and Sb are the re excess minority carriers recombination at the junction [17, 19, 29] and the rear [20, 21, 23] of the base.

### 3. Results and Discussions

#### 3.1 Photocurrent density

The AC photocurrent density at the junction is obtained from minority carriers density in the base and is given by the following expression:

$$J_{ph}(Sf, z, H, \omega, \alpha(\lambda)) = qD(\omega) \left. \frac{\partial \delta(x, z, H, Sf, \omega, \alpha(\lambda))}{\partial x} \right|_{x=0} \tag{13}$$

Where q is the elementary electron charge.

#### Influence of modulation frequency on photocurrent density

Figure 3 shows the influence of frequency on the photocurrent density, which is very marked at large (Sf) values corresponding to the solar cell in short circuit condition illuminated with strong penetration absorption coefficient ( $\alpha = 6.2 \text{ cm}^{-1}$ ).

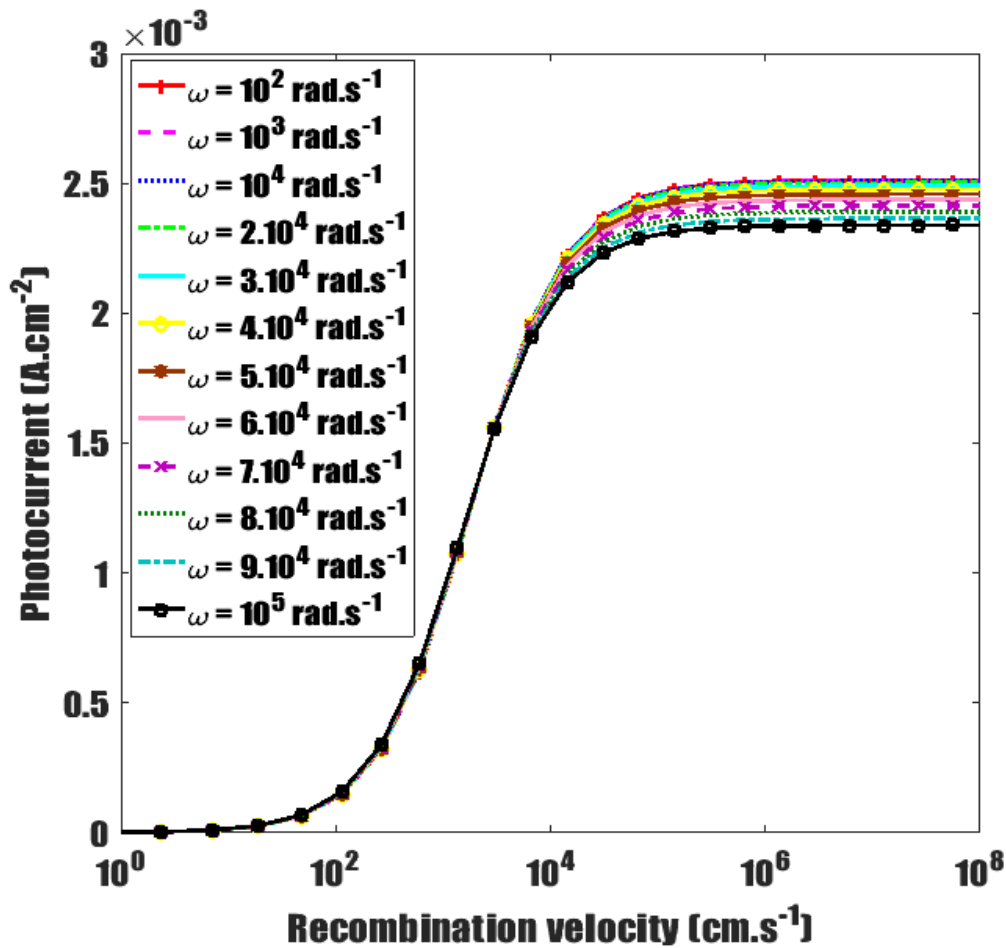


Figure 3: Profile of the photocurrent module density as a function of the recombination velocity at the junction, for different frequency values ( $D_0 = 35\text{cm}$ ;  $\alpha = 6.2 \text{ cm}^{-1}$ ;  $z = 0,017 \text{ cm}$ ;  $H=0,02 \text{ cm}$ )



### 3.2 Base thickness optimization

The representation of photocurrent density (Figure. 3) according to the junction recombination velocity of minority carriers shows that, for very large Sf, a bearing sets up and corresponds to the short-circuit current density (Jphsc). So, in this junction recombination velocity interval [16, 30], we can write:

$$\left. \frac{\partial J_{ph}(Sf, z, H, \lambda, \omega)}{\partial Sf} \right|_{Sf \geq 10^5 \text{ cm.s}^{-1}} = 0 \tag{14}$$

The solution of equation (14) leads to the expressions of ac recombination velocity in the back surface, given by equations (15) and (16):

$$Sb1(\omega, H) = -\frac{D(\omega)}{L(\omega)} \cdot th\left(\frac{H}{L(\omega)}\right) \tag{15}$$

$$Sb2(\omega, H) = -\frac{D(\omega)sh\left(\frac{H}{L(\omega)}\right)}{L(\omega)\left(ch\left(\frac{H}{L(\omega)}\right) - 1\right)} \tag{16}$$

Figure 4 representation is the profile of the two expressions of back surface recombination velocity versus thickness of the base of the solar cell, in order to determine the optimum thickness of the base, abscissa of the intercept point, for given frequency.

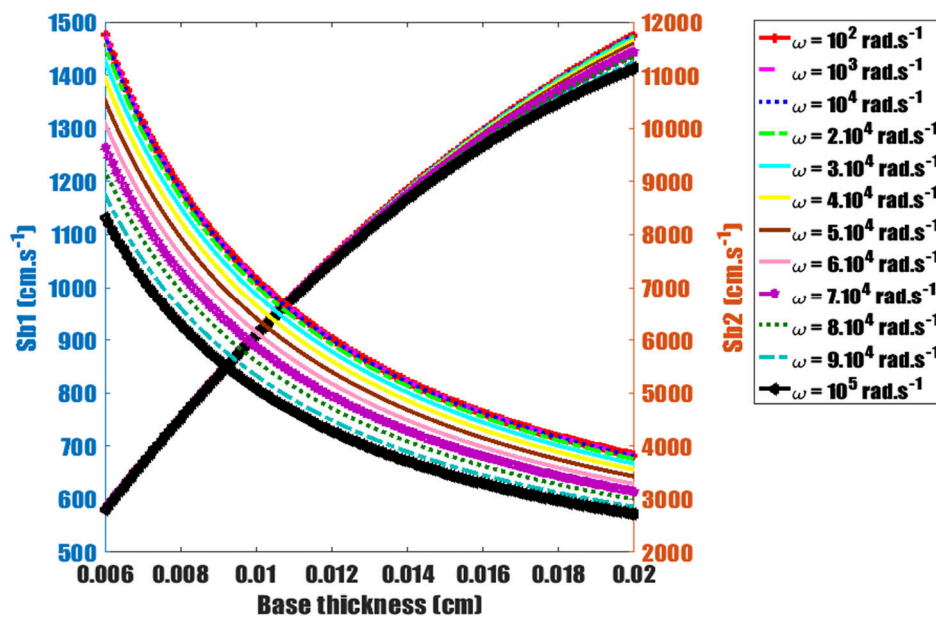


Figure 4: Sb1 and Sb2 versus depth in the base for different frequency

The intercept point of the two curves leads to the optimum thickness of the base for each modulation frequency. Table 1 shows the results obtained. Figures 4 and 5 produce the (Hopt) representation as a function of the modulation frequency and the dynamic diffusion coefficient respectively.

$\omega(\text{rad.s}^{-1})$	$10^2$	$10^3$	$10^4$	$2.10^4$	$3.10^4$	$4.10^4$	$5.10^4$	$6.10^4$	$7.10^4$	$8.10^4$	$9.10^4$	$10^5$
D (cm <sup>2</sup> /s)	35	34.99 65	34.65 35	33.65 38	32.11 01	30.17 24	28.00 00	25.73 53	23.48 99	21.34 15	19.33 70	17.50 00
H <sub>opt</sub> (cm)	0.01 07	0.010 7	0.010 7	0.010 6	0.010 5	0.010 4	0.010 2	0.01	0.009 8	0.009 6	0.009 4	0.009 2

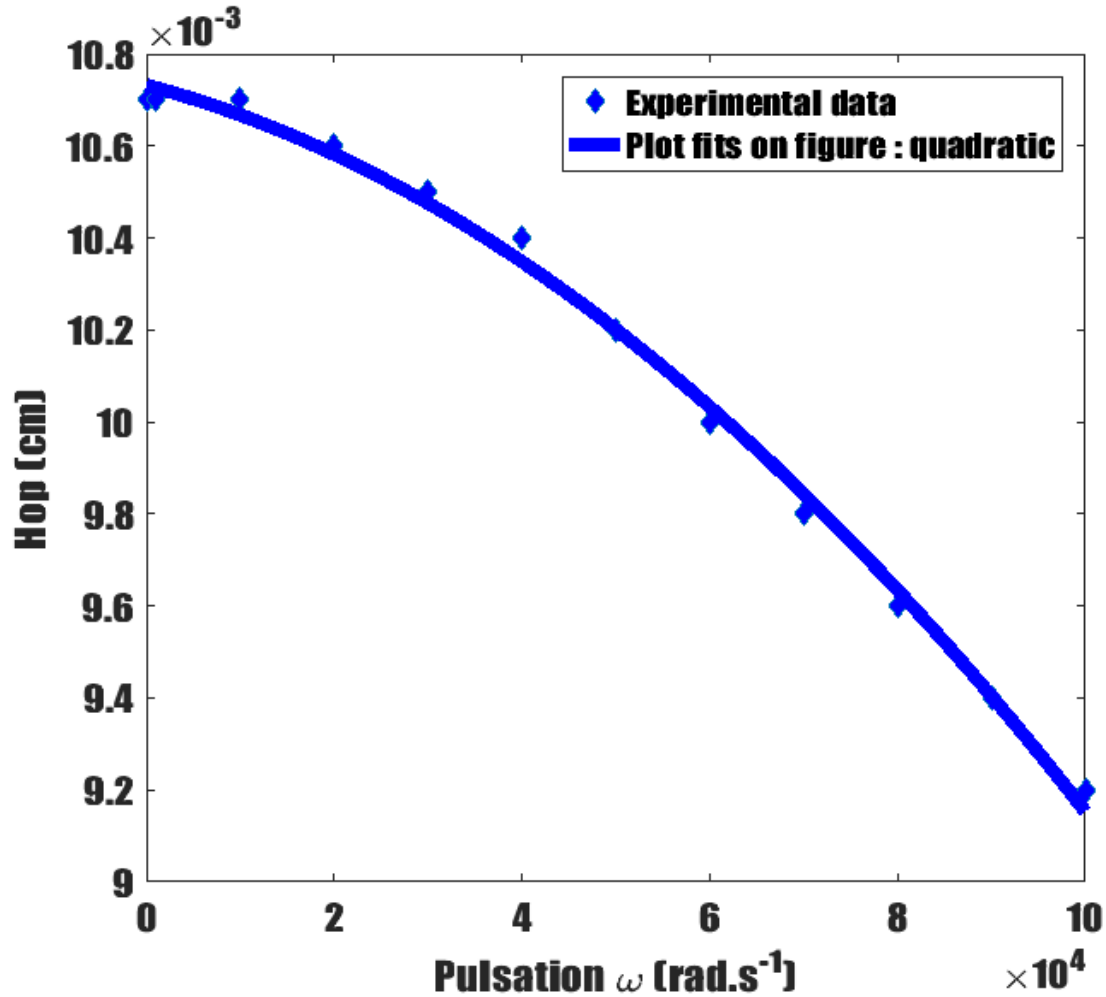


Figure 5: Optimum thickness versus frequency

Equation (17) gives the mathematical correlation of the optimum thickness with the modulation frequency. Hopt is a decreasing function of the modulation frequency, expressed as:

$$Hopt(cm) = -10^{-13} \times \omega^2 - 5.5 \cdot 10^{-9} \times \omega + 0.011 \tag{17}$$

Thus, it appears that low frequencies require significant base thicknesses, while at high frequencies, the optimum thickness remains thin [31, 32, 33; 34].

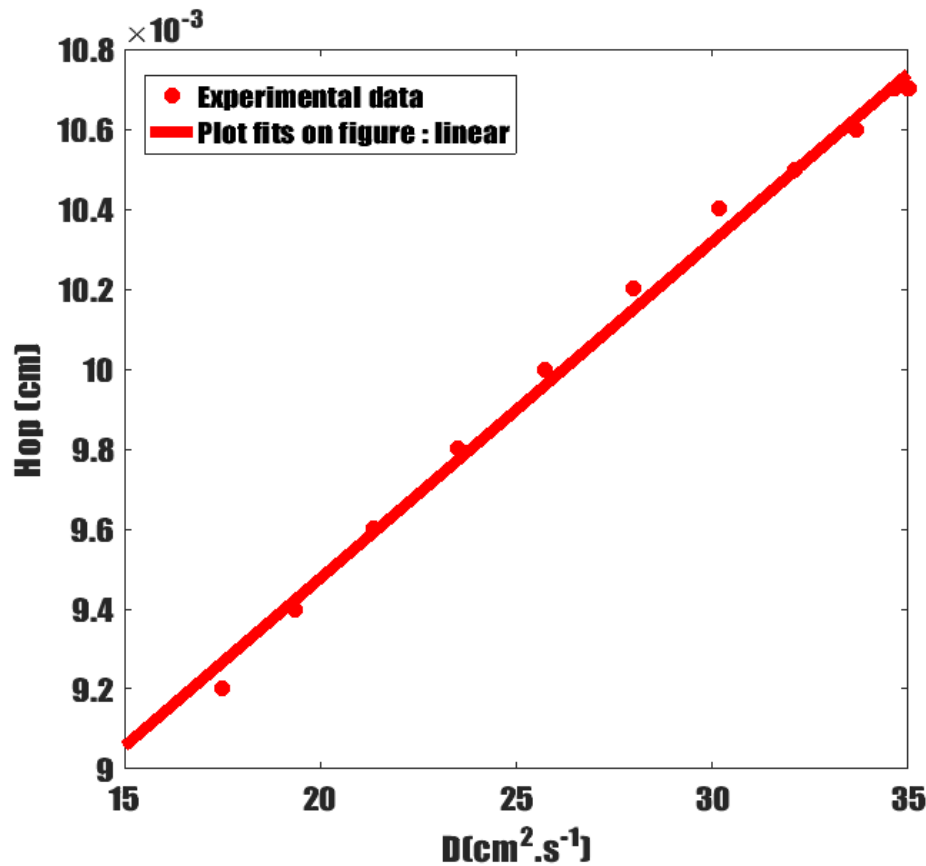


Figure 6: Optimum thickness versus diffusion coefficient

The equation (18) is an increasing  $H_{opt}$  function with the dynamic diffusion coefficient. These latter (Eq.5) is decreasing with the modulation frequency.

$$H_{opt}(cm) = 8.4 \cdot 10^{-5} \times D + 0.0078 \quad (18)$$

Various research works have produced results on the optimum thickness of the base of the silicon solar cell, which actually increases with the diffusion coefficient of the minority carriers in the base subjected to physical conditions, such as, monochromatic illumination [35], magnetic field [36, 37], temperature [38], irradiation flux by charged particles [39], variation in doping rate [34, 40], as well as the possible combination of these different factors [41, 42, 43].

The optimum thickness, on the other hand, decreases with these factors and makes it possible to model the conditions allowing the saving of matter in the development of the solar cell [44, 45, 46].

#### 4 Conclusion

The technique of intersecting the curves of the recombination rates of the minority carriers in the base was applied on silicon solar cell, with multi vertical junctions connected in series and placed under illumination in frequency modulation.

The optimum thickness of the base decreases with the modulated frequency, while it increases with minority carriers' dynamic diffusion coefficient.

The modeling equations of the optimum thickness of the base indicate how to achieve the saving of silicon material in the development of the solar cell.



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